

Influence of working pressure and lasing energy of Al plasma in laser-induced breakdown spectroscopy

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Abstract

Aluminum plasma was generated by the irradiation of the target with Nd: YAG laser operated at a wavelength of 1064 nm. The effect of laser power density and the working pressure on spectral lines generating by laser ablation, were detected by using optical spectroscopy. The electron density was measured using the Stark broadening of aluminum lines and the electron temperature by Boltzmann plot method it is one of the methods that are used. The electron temperature T_e , electron density n_e , plasma frequency ω_p and Debye length λ_D increased with increasing the laser peak power. The electron temperature decrease with increasing gas pressure.

Key words

LIBS, plasma characteristics, OES, electron temperature, electron density.

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تأثير ضغط التشغيل وطاقة الليزر على بلازما الالمنيوم باستخدام التحليل الطيفي المستحث بالليزر

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الخلاصة

توليد بلازما الالمنيوم بتشعيع هدف من معدن الالمنيوم بليزر نديميوم-ياك ذو طول موجي 1064 نانومتر. تم تحديد تأثير التغيير في طاقات الليزر وضغوط العمل على خط الطيف المتولد بواسطة القلع بالليزر، تحديد معاملات البلازما باستخدام التحليل الطيفي، كثافة الالكتران تم قياسها باستخدام طريقة توسيع ستارك لخطوط الالمنيوم ودرجة حرارة الالكتران باستخدام منحنيات طريقة بولتزمان، لوحظ في النتائج ان درجة حرارة الالكتران، كثافته، تردد البلازما وطول ديبي تزداد بزيادة طاقة الليزر. درجة الحرارة تقل بزيادة الضغط للغاز.

Introduction

LIBS is an atomic emission *spectrum analysis* technique. Which is a high pulsed laser power density which is localized onto a small target volume leading to a breakdown of analytical into ions and free electrons, resulting a plasma identify by both continuum and atomic emission [1]. The parameters of the Laser-generated plasma is progress quickly and are

strongly related on irradiation conditions like Intensity of laser incident on a surface of the target and pulse laser duration, the wavelength of the laser, the spot size of radiance, gas ambient installation, and ambient pressure, which are temporary in nature [2, 3]. Optical Emission Spectroscopy (OES) has been used for years to determine the plasma parameters such as electron density

(n_e) and electron temperature (T_e), Debye length (λ_D) and plasma frequency (f_D) [4]. Plasma diagnostic method can be done through an account of the plasma electron temperature (T_e) and density (n_e). The strength of different distribution functions describing the plasma state determines by the temperature, while the state of thermo-dynamical equilibrium of the plasma specifies by the electron density [5]. For the calculation of electron temperature the Boltzmann plot method it is one of the simplest methods that are used [6, 7].

$$\ln(\lambda_{mn} I_{mn} / g_m A_{mn}) = -E_m / K T_e + \ln(N(T) / U(T)) \quad (1)$$

where: λ_{mn} is wavelength of the spectral line, I_{mn} intensity of the spectral line. A_{mn} is the transition probability of upper (m) and lower (n) energy states. E_m and g_m are the upper states (m) energy and statistical weight T_e is the electron temperature, K is the Boltzmann constant, $N(T)$ and $U(T)$ are the total number density and partition function.

The Stark broadening effect was used to calculate the electron density requires a line which is free from self-absorption [8]:

$$n_e = \left[\frac{\Delta\lambda}{2\omega_s} \right] N_r \quad (2)$$

ω_s : is the theoretical line full width Stark broadening parameter, calculated at the same reference electron density $N_r \approx 10^{17} \text{ cm}^{-3}$. Plasma frequency can be given as [9]:

$$\omega_p = (n_e e^2 / m_e \varepsilon_0)^{1/2} \quad (3)$$

where ε_0 is the Permittivity of free

space, m_e is the electron mass, n_e is the electron number density and e is electron charge.

Debye's length can be calculated by the formula [9]:

$$\lambda_D = \left[\frac{\varepsilon_0 K_B T_e}{n_e e^2} \right]^{1/2} \cong 7.43 \times 10^2 \left(\frac{T_e (\text{eV})}{n_e} \right)^{1/2} \quad (4)$$

where: k_B is Boltzmann's constant, T_e is the plasma temperature, e is the electron charge and n_e is the electron density.

Experimental part

Fig.1 illustrates the experimental setup of LIBS system used for the detection of the spectral lines of the laser-generated Al plasma, target was bombarded by Nd: YAG pulse laser (9 ns duration, 6 Hz repetition frequency, and fundamental wavelength of 1064 nm) and laser pulse energy ranging from 100 mW to 400 mW on target surface, at an angle of 45° . To focus the laser on the target surface we used the convex lens with a focal length equal to 10 cm. The circularly shaped sheet of the aluminum target with diameter 3 cm is placed inside a vacuum chamber. The chamber was filled up with Argon gas at a pressure (0.2 and 0.4 Torr). The vacuum chamber made of a cylindrical stainless steel tube. The two ends closed by Pyrex windows, by two stainless steel flanges, and with small quartz window fixed in its center, that allows the laser pulse to shoot the Al target. Two small pipes connected to pumping systems, while the other was used to deliver the argon gas with purity (99.9 %).

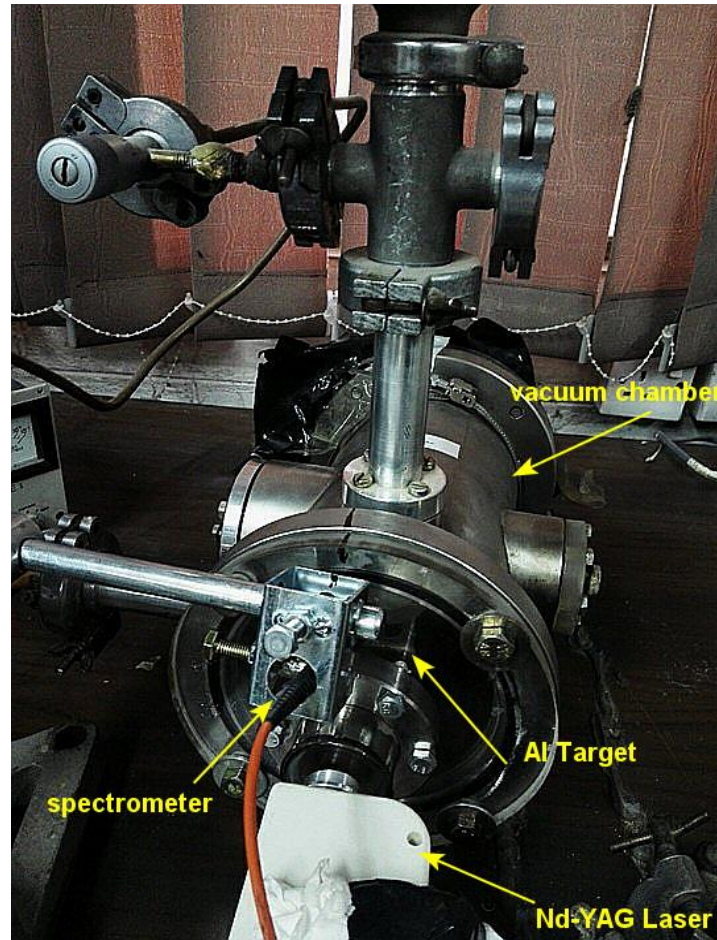


Fig.1: Experimental set up of the LIBS system.

Results and discussions

The Al plasma was produced by the laser interaction with Al target using Q-switched Nd: YAG in a vacuum. A spectrum consists of a number of characteristic spectral lines of particular atoms and ions. The increase in the pulse laser power density will increase its absorption in the plasma resulting in more ablation which results in the increase of the spectral line intensities of the plasma emission. The value of T_e is obtained using the Boltzmann plot method (Eq.(1)), this requires peaks that

originated from the same atomic species and the same ionization stage with data from NIST site, where the electron temperature equal to the invert of the slope of the fitting line (the slope of fitted line equals to $(-1/k_B T)$). Fig. 2 and 3 show a Boltzmann plot by plotting a graph between $\ln\left(\frac{I_{ji}\lambda_{ji}}{hc g_j A_{ji}}\right)$ as a function of E_k . The electron temperature of Al plasma increases with increasing laser power from (100 to 400 mW), electron temperature decrease with increasing working pressure.

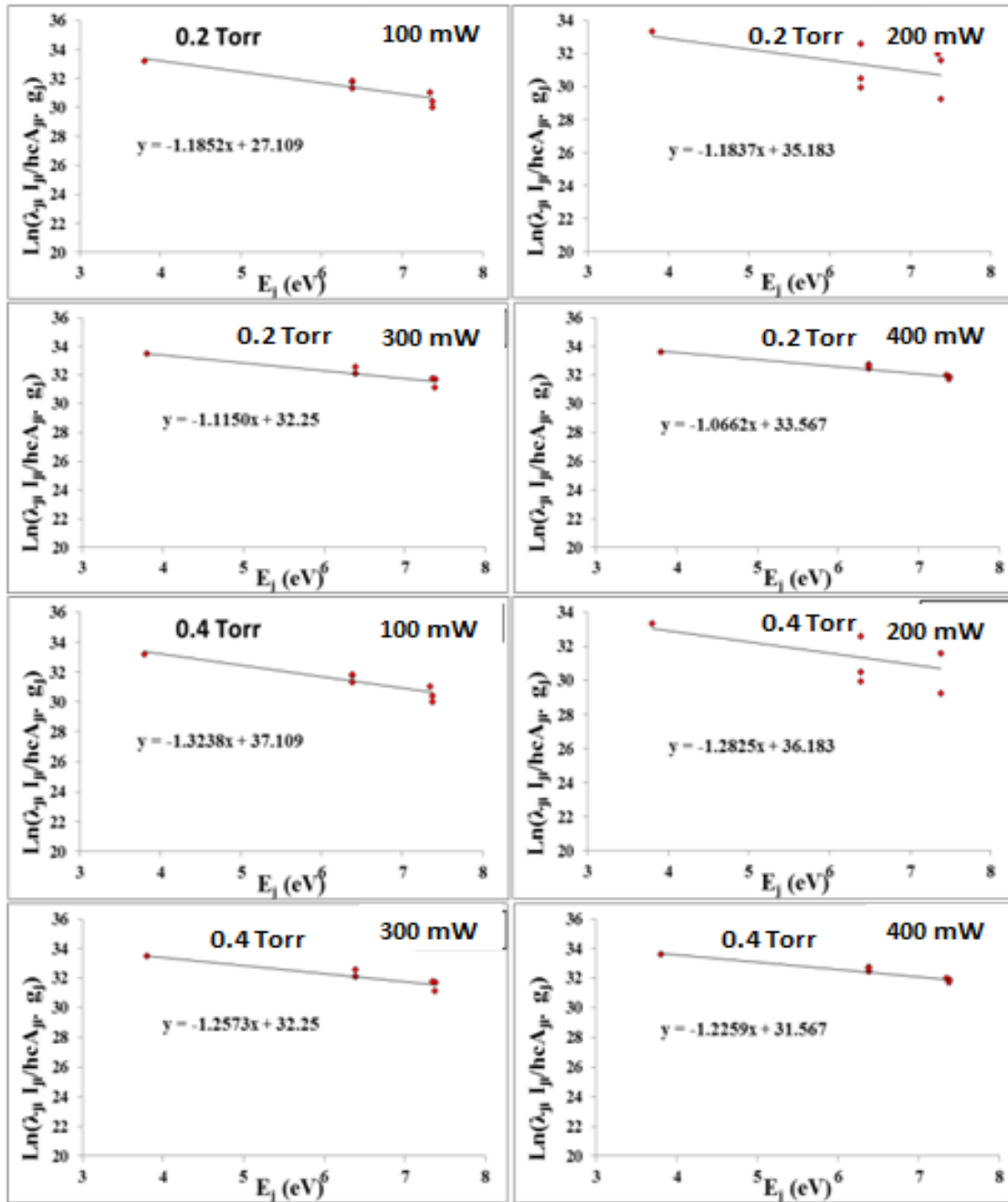


Fig. 2: Boltzmann plot for Al target with different laser peak power in 0.2 and 0.4 Torr at 1064 nm.

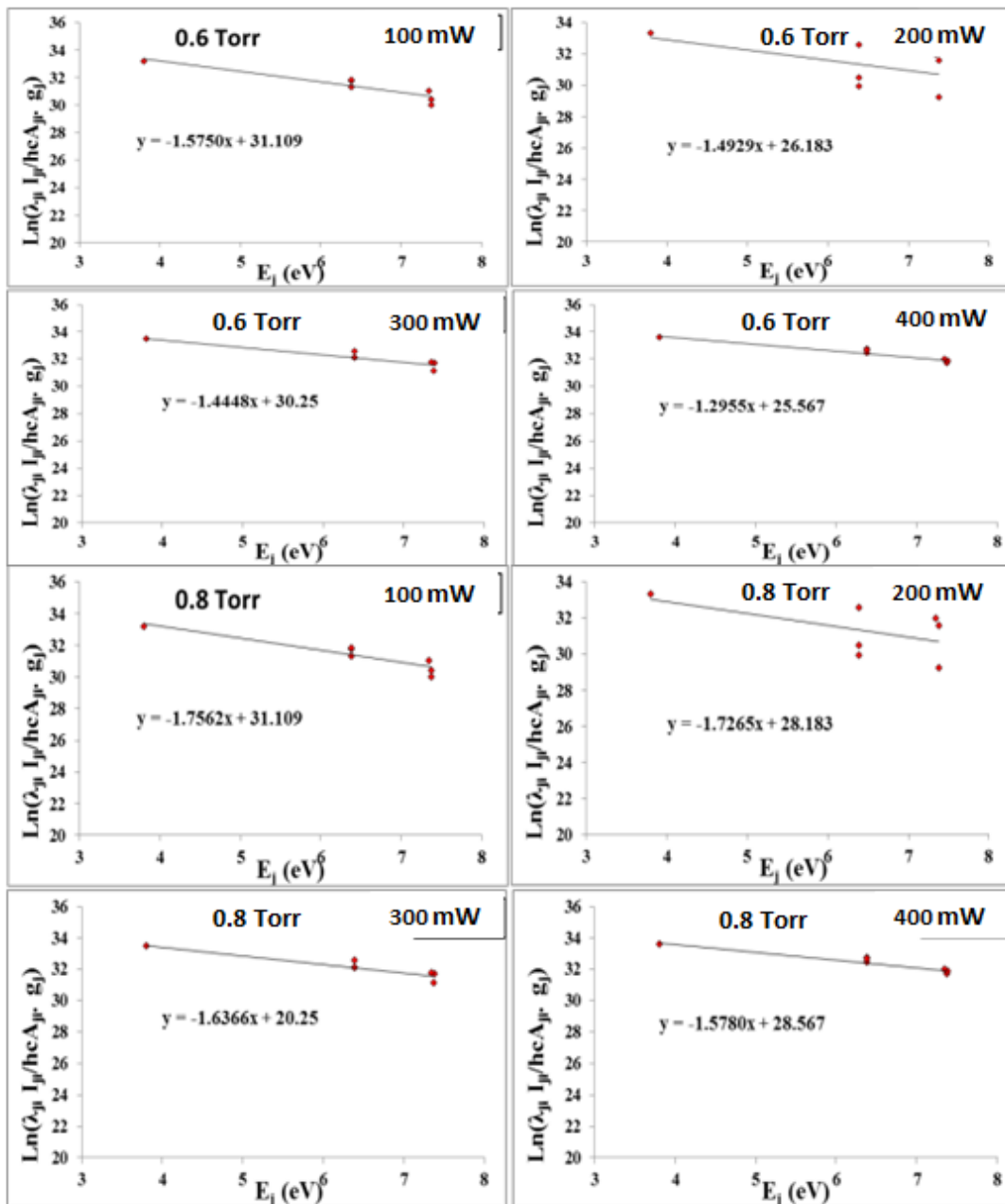


Fig. 3: Boltzmann plot for Al target with different laser peak power in 0.6 and 0.8 Torr at 1064 nm.

Fig. 4 shows the variation of T_e and n_e with laser power density. The results of these figures illustrated that the increase of laser power density lead to an increase of electron temperature and the electron density, this is due to the

absorption of the laser photon by the plasma [10]. The increasing of n_e with increasing of laser power density and working pressure attributed to the increase of electron collisions with the ambient gas.

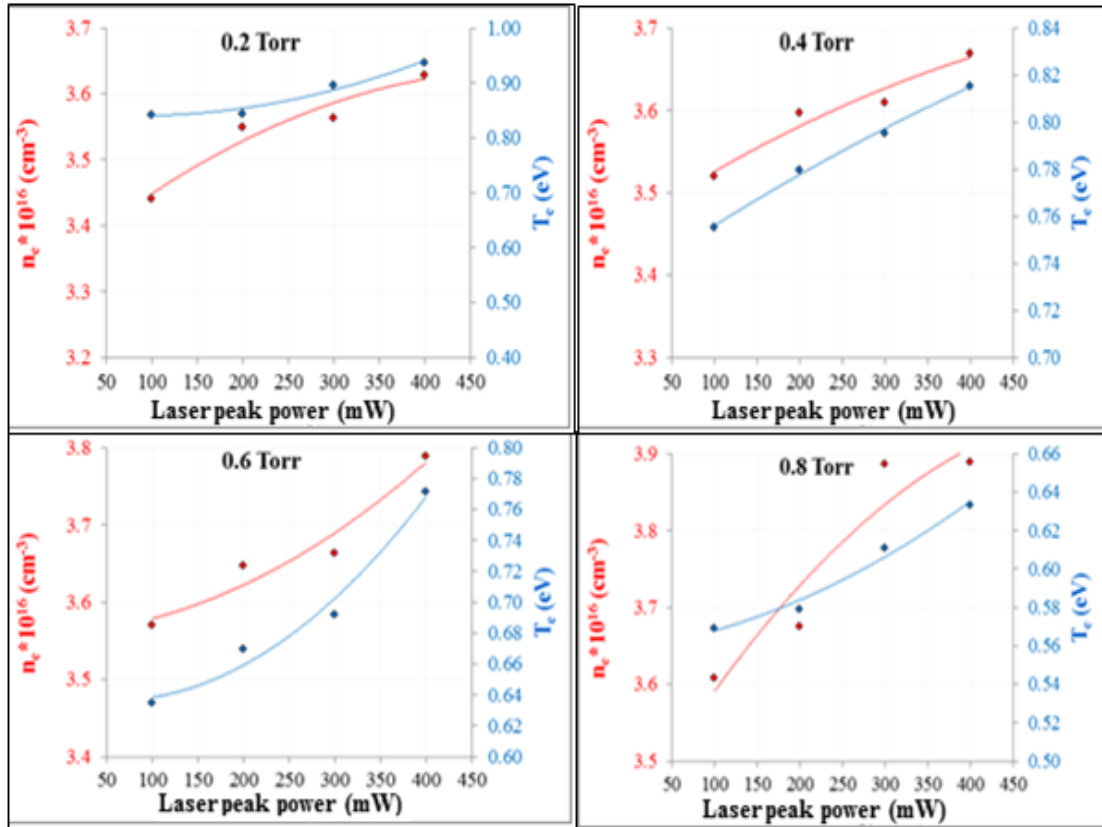


Fig.4 : The variation of T_e and n_e versus the laser peak power.

Table 1 shows the values of electron temperature (T_e), electron density (n_e), Debye length (λ_D) by using Eq. (4), plasma frequency (ω_p) using eq.3 for Al targets at different laser pulse energies (100 to 400) mJ with different working pressures (0.2 to 0.8) Torr. The values of plasma frequency increase because of the

dependence of plasma frequency on the density of electrons (n_e)^{1/2}, We can also notice that the plasma frequency increases with the increase of the laser energies. Also the, Debye length increase because Debye length depends on plasma temperature and plasma density [11].

Table 1: Plasma parameters for Al target at (1064 nm) in vacuum with different laser energies and different working pressure.

0.2 Torr	Laser peak power (mW)	Te (eV)	$n_e \cdot 10^{16} \text{ (cm}^{-3}\text{)}$	$\omega p \cdot 10^{12} \text{ (Hz)}$	$\lambda_D \cdot 10^{-5} \text{ (cm)}$
	100	0.7837	3.1383	1.5908	0.3713
	200	0.8007	3.2411	1.6166	0.3693
	300	0.8021	3.3319	1.6391	0.3645
	400	0.8473	3.4358	1.6645	0.3697
0.4 Torr	Laser peak power (mW)	Te (eV)	$n_e \cdot 10^{16} \text{ (cm}^{-3}\text{)}$	$\omega p \cdot 10^{12} \text{ (Hz)}$	$\lambda_D \cdot 10^{-5} \text{ (cm)}$
	100	0.7041	3.1988	1.6060	0.3486
	200	0.7360	3.2980	1.6308	0.3510
	300	0.7509	3.4251	1.6619	0.3479
	400	0.7629	3.5261	1.6862	0.3456
0.6 Torr	Laser peak power (mW)	Te (eV)	$n_e \cdot 10^{16} \text{ (cm}^{-3}\text{)}$	$\omega p \cdot 10^{12} \text{ (Hz)}$	$\lambda_D \cdot 10^{-5} \text{ (cm)}$
	100	0.5509	3.2455	1.6177	0.3061
	200	0.6017	3.3043	1.6323	0.3170
	300	0.6403	3.4521	1.6684	0.3200
	400	0.6912	3.5594	1.6942	0.3274
0.8 Torr	Laser peak power (mW)	Te (eV)	$n_e \cdot 10^{16} \text{ (cm}^{-3}\text{)}$	$\omega p \cdot 10^{12} \text{ (Hz)}$	$\lambda_D \cdot 10^{-5} \text{ (cm)}$
	100	0.5312	3.3776	1.6503	0.2946
	200	0.5593	3.3909	1.6536	0.3017
	300	0.5635	3.5630	1.6950	0.2954
	400	0.5701	3.6307	1.7110	0.2944

Conclusions

The plasma parameters such as temperature and number density are found to increase with the laser irradiance. It is inferred that at first stage the laser vapor interaction is largely due to the inverse bremsstrahlung (IB) process. And the electron temperature decrease with increasing working pressure from 0.2 Torr to 0.8 Torr, The increasing of inelastic collisions of an electron with Ar atoms with increasing of gas

pressure is responsible for decreasing of T_e with increasing of pressure, the electron density increase with increasing the working pressure.

References

- [1] J. Feng, Z. Wang, Z. Li, W. Ni, Spectrochim. Acta - Part B At. Spectrosc., 65, 7 (2010) 549-556.
- [2] V. N. Rai, H. Zhang, F. Y. Yueh, J. P. Singh, A. Kumar, Appl. Opt., 42, 18 (2003) 3662-3669.
- [3] S. S. Harilal, C. V. Bindhu, M. S.

Tillack, F. Najmabadi, A. C. Gaeris, J. Appl. Phys., 93, 5 (2003) 2380-2388.

[4] Ulf Saalman and Jan-Michael Rost, Phys. Rev. Lett., 89, 14 (2002) 143401_1-143404.

[5] A. M. El Sherbini, World J. Nano Sci. Eng., 2, December (2012) 206-212.

[6] J. A. Aguilera and C. Aragón, Spectrochim. Acta-Part B At. Spectrosc., 59, 12 (2004) 1861-1876.

[7] S. Waheed, Shazia Bashir, Asadullah Dawood, Safia Anjum, Mahreen Akram, Asma Hayat, Saba Amin, Ali Zaheer, International Journal for Light and Electron Optics,

International Journal for Light and Electron Optics, 140 (2017) 536-544.

[8] A. M. El Sherbini, Opt. Photonics J., 2, 4 (2012) 278-285.

[9] R. K. Tyagi, R. S. Pandey, A. Kumar, K. K. Srivastava, Int. J. Eng. Sci. Technol., 3, 8 (2011) 168-176.

[10] M. V. Allmen and A. Blatter, Springer-Verlag, Berlin, 2nd ed, (1995).

[11] Bhatti, K. A., Khaleeq Rahman, M. Rafique, M. S. Shahzad, M. I. Latif, A. Parveen, Vacuum, 82, 11 (2008) 1157-1161.